

Table 3. Tetrahedral position angles, γ , and the ratios of octahedral–dodecahedral shared edge (O \cdots O v) to dodecahedral–dodecahedral shared edge (O $^v\cdots$ O ix) distances (Å) in $M_3^{2+}Al_2Si_3O_{12}$ garnets

	γ^*	(O \cdots O v)/(O $^v\cdots$ O ix)	Ratio	Reference
Co(α)	28.07	2.639 (3)/2.705 (4)	0.976	Ohashi <i>et al.</i> (1981)
Co(β)	27.75	2.625 (1)/2.713 (1)	0.968	Present study
Mg	27.42	2.617 (3)/2.709 (3)	0.966	Novak & Gibbs (1971)
Fe	26.69	2.655 (2)/2.768 (3)	0.959	Novak & Gibbs (1971)
Mn	26.08	2.678 (2)/2.824 (3)	0.948	Novak & Gibbs (1971)
Ca, Mn	25.25	2.692 (2)†/2.879 (2)†	0.935	Novak & Gibbs (1971)
Ca	24.50	2.756 (2)/2.973 (2)	0.927	Novak & Gibbs (1971)
Cd	24.38	2.739 (8)/2.970 (10)	0.922	Ohashi <i>et al.</i> (1985)

* Born & Zemann (1964).

† Calculated by present authors.

The SDP program system (B. A. Frenz & Associates Inc., 1982) was employed.

Lists of structure factors and anisotropic displacement parameters have been deposited with the IUCr (Reference: OH1074). Copies may be obtained through The Managing Editor, International Union of Crystallography, 5 Abbey Square, Chester CH1 2HU, England.

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Redetermination of Sodium Cerium(III) Sulfate Monohydrate, NaCe(SO $_4$) $_2$ ·H $_2$ O

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Abstract

The structure of sodium cerium sulfate monohydrate, NaCe^{III}(SO $_4$) $_2$ ·H $_2$ O, comprises distorted NaO $_6$ octahedra, CeO $_9$ polyhedra in the form of distorted tricapped trigonal prisms, and slightly irregular tetrahedral sulfate ions. This structure is isomorphic with that of NaLa^{III}(SO $_4$) $_2$ ·H $_2$ O. All bonds fall within normal lim-

its. The sulfate ion manifested rigid-body behavior but neither of the cation complexes did. As in the previous structural analysis of the La analog [Blackburn & Gerkin (1994). *Acta Cryst.* **C50**, 835–838], the water molecule is modelled as having O-atom disorder but with an ordered H atom. Hydrogen bonds involve only sulfate O atoms as acceptors. The anisotropic atomic displacement parameters are found to be in good agreement with the corresponding parameters determined for the La analog, but in poor agreement with those previously reported for the title salt by Lindgren [*Acta Chem. Scand. Ser. A*, (1977), **31**, 591–594]. Taken together, the data establish that the lack of correction for absorption in the previous study of the Ce salt, rather than some intrinsic property of the Ce salt, is responsible for the displacement parameter discrepancies. The absolute structure has been determined.

Comment

This study of NaCe^{III}(SO $_4$) $_2$ ·H $_2$ O was performed, following our study of the La^{III} analog (Blackburn & Gerkin, 1994), to allow direct comparisons with earlier results for the title salt (Lindgren, 1977). Of particular interest was the characterization of the anisotropic atomic displacement parameters since the corresponding values found in our study of the La analog differed notably from those reported for the Ce salt by Lindgren (1977). Moreover, the structural characterization of the water molecule was also of particular interest since in the La analog the water O atom was best modelled as disordered equally over two sites, but Lindgren's analysis did not ascribe such disorder to the Ce salt.

The experimental crystal belonged to the enantiomeric space group $P3_221$ (see *Experimental*), whereas our experimental crystal of the La analog and Lindgren's Ce crystal belonged to $P3_121$.

This salt is isomorphic with the La analog except for its handedness. The atom numbering corresponds with that chosen for the La analog, and structural descriptions given for the La analog (Blackburn & Gerkin, 1994) apply here with only minor changes.

In the distorted octahedral NaO $_6$ complex, the Na—O distances range from 2.447 (4) to 2.538 (3) Å [average 2.49 (4) Å] and the O—Na—O angles range from 71.7 (1) to 154.2 (2) $^\circ$ (Table 2). The next nearest atoms to Na, at distances of 2.849 (3) Å, are a symmetry-related pair of O(4) atoms. In the CeO $_9$ complex, the Ce—O distances range from 2.477 (3) to 2.582 (3) Å [average 2.53 (3) Å] while the O—Ce—O angles range from 54.4 (1) to 145.7 (1) $^\circ$ (Table 2). These values are in good agreement with our previous values for the La analog and also with those of Lindgren.

The sulfate ion is subject to no symmetry constraints and forms a slightly irregular tetrahedron. The observed S—O distances range from 1.465 (3) to 1.481 (3) Å [average 1.473 (6) Å] and the O—S—O angles range

from 105.7(2) to 112.0(2)° (Table 2). These values very closely match the corresponding values for the La analog. Lindgren's values for S—O distances ranged from 1.458(13) to 1.490(11) Å [average 1.479(13) Å] and the O—S—O angles ranged from 106.4(6) to 111.6(7)°. Thus, while the angular ranges and the average S—O bond distances are in good agreement, the range of S—O bond distances reported here for the Ce salt is only approximately one half of that reported earlier by Lindgren (the difference, however, is not statistically significant). As in the La analog, the hydrogen bonding in this salt involves only a single O donor atom, O(5), and only sulfate O atoms as acceptors. Again, only O(2) and O(4) are involved in substantial hydrogen bonds (Table 3).

Notable discrepancies between the present results and the earlier results of Lindgren are found in the anisotropic displacement parameters. Any numerical comparisons are rendered moot since Lindgren's U_{ij} values were found to yield several negative principal values for the displacement ellipsoids. Since our principal values for the title salt were uniformly positive, and are in good agreement with those of the La analog, the problem with Lindgren's analysis almost certainly arose from the lack of correction for absorption, rather than from some intrinsic property of his sample of the Ce salt.

The presence in the sample from gel crystal growth of single crystals of the title salt belonging to the enantiomeric space group $P3_121$ was confirmed.

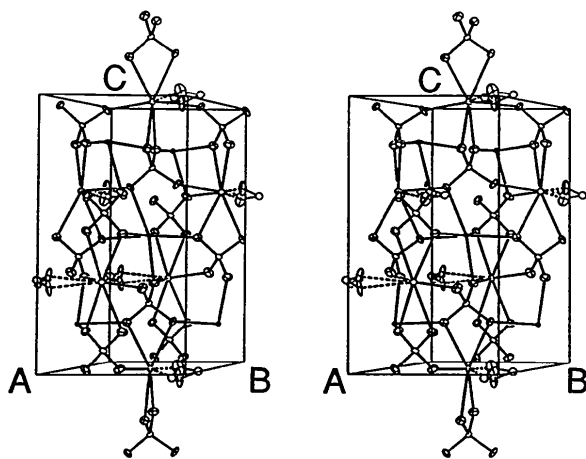


Fig. 1. An ORTEP (Johnson, 1976) stereoview of a unit cell of sodium cerium(III) sulfate monohydrate at 296 K. Displacement ellipsoids have been drawn at 50% probability for all atoms except H, for which they have been set artificially small. The Na ions have been filled for contrast. Both positions of the disordered water O atom are shown. Hydrogen bonds have been omitted for clarity.

Experimental

Single crystals of NaCe^{III}(SO₄)₂·H₂O were prepared by gel-diffusion growth as described previously for the La^{III} analog (Blackburn & Gerkin, 1994). The morphology of the crystals

was strikingly similar to that of the La analog: the crystals have six pentagonal side faces {010} with three capping pentagonal faces {101} (space group $P3_121$) or $\{\bar{1}01\}$ (space group $P3_221$) at each end. The experimental sample was mounted with epoxy cement on a glass fiber.

Crystal data

NaCe(SO₄)₂·H₂O

$M_r = 373.24$

Trigonal

$P3_221$

$a = 7.019(1)$ Å

$c = 12.925(2)$ Å

$V = 551.4(2)$ Å³

$Z = 3$

$D_x = 3.372$ Mg m⁻³

Mo $K\alpha$ radiation

$\lambda = 0.71073$ Å

Cell parameters from 25 reflections

$\theta = 16.3\text{--}17.4^\circ$

$\mu = 6.92$ mm⁻¹

$T = 296$ K

Capped hexagonal prism

$0.38 \times 0.13 \times 0.13$ mm

Colorless

Data collection

AFC-5S diffractometer

ω scans

Absorption correction:

analytical

$T_{\min} = 0.315$, $T_{\max} =$

0.489

1877 measured reflections

848 independent reflections

831 observed reflections

$[I > 3\sigma(I)]$

$R_{\text{int}} = 0.017$

$\theta_{\max} = 27.5^\circ$

$h = 0 \rightarrow 7$; $-7 \rightarrow 0$

$k = 0 \rightarrow 7$; $-7 \rightarrow -1$

$l = -16 \rightarrow 16$; $-16 \rightarrow 16$

6 standard reflections

monitored every 150

reflections

intensity variation: $\pm 1.8\%$

(average maximum

relative intensity)

Refinement

Refinement on F

$R = 0.017$

$wR = 0.030$

$S = 1.72$

831 reflections

60 parameters

H-atom parameters not

refined (see below)

$w = 1/\sigma^2(F)$

$(\Delta/\sigma)_{\max} < 0.01$

$\Delta\rho_{\max} = 1.63$; next

highest = 0.56 e Å⁻³

$\Delta\rho_{\min} = -0.56$ e Å⁻³

Extinction correction: none

Atomic scattering factors

from Stewart, Davidson &

Simpson (1965) (H) and

Cromer & Waber (1974)

(O, Na, S, Ce)

Table 1. Fractional atomic coordinates and equivalent isotropic displacement parameters (Å²)

$$U_{\text{eq}} = (1/3)\sum_i\sum_j U_{ij}a_i^*a_j^*a_i\cdot a_j$$

	x	y	z	U_{eq}
Ce	0.43376 (5)	0	2/3	0.00955 (10)
S	0.4528 (2)	0.0115 (2)	0.91277 (7)	0.0088 (2)
Na	0.4656 (3)	0	1/6	0.0070 (5)
O(1)	0.4972 (6)	-0.1215 (6)	0.9856 (2)	0.0177 (8)
O(2)	0.2599 (5)	-0.1308 (5)	0.8474 (2)	0.0157 (7)
O(3)	0.6372 (6)	0.1339 (6)	0.8404 (2)	0.0190 (8)
O(4)	0.4100 (6)	0.1627 (6)	0.9738 (2)	0.0188 (8)
O(5)†	0.0716 (11)	-0.0076 (18)	0.6450 (6)	0.0305 (16)

† Occupancy factor 0.5.

Table 2. Selected geometric parameters (Å, °)

Ce—S	3.183 (1)	S—O(1)	1.465 (3)
Ce—S'	3.183 (1)	S—O(2)	1.481 (3)
Ce—O(1 ^u)	2.477 (3)	S—O(3)	1.475 (4)

Ce—O(1 ⁱⁱⁱ)	2.477 (3)	S—O(4)	1.469 (3)
Ce—O(2)	2.582 (3)	Na—O(1 ^{vi})	2.538 (3)
Ce—O(2')	2.582 (3)	Na—O(1')	2.538 (3)
Ce—O(3)	2.573 (3)	Na—O(2 ^{viii})	2.493 (4)
Ce—O(3')	2.573 (3)	Na—O(2 ^{viiii})	2.493 (4)
Ce—O(4 ^{iv})	2.496 (3)	Na—O(3 ^{ix})	2.447 (4)
Ce—O(4')	2.496 (3)	Na—O(3 ^x)	2.447 (4)
Ce—O(5)	2.530 (6)	O(5)—H	0.97
Ce—O(5')	2.530 (6)	O(5)—H ⁱ	0.97
O(1 ⁱⁱ)—Ce—O(1 ⁱⁱⁱ)	139.7 (2)	O(1)—S—O(2)	110.6 (2)
O(1 ⁱⁱ)—Ce—O(2)	71.22 (10)	O(1)—S—O(3)	112.0 (2)
O(1 ⁱⁱ)—Ce—O(3)	119.3 (1)	O(1)—S—O(4)	107.5 (2)
O(1 ⁱⁱ)—Ce—O(4 ^{iv})	145.5 (1)	O(2)—S—O(3)	105.7 (2)
O(1 ⁱⁱ)—Ce—O(5)	68.4 (3)	O(2)—S—O(4)	110.2 (2)
O(1 ⁱⁱ)—Ce—O(5')	71.5 (3)	O(3)—S—O(4)	110.9 (2)
O(2)—Ce—O(2')	145.7 (1)	O(1 ^{vi})—Na—O(1')	150.4 (2)
O(2)—Ce—O(4 ^{iv})	122.68 (9)	O(1 ^{vi})—Na—O(2 ^{viii})	71.7 (1)
O(2)—Ce—O(5)	78.5 (2)	O(1')—Na—O(2 ^{viii})	85.4 (1)
O(2)—Ce—O(3)	54.4 (1)	O(1 ^{vi})—Na—O(3 ^{ix})	79.3 (1)
O(3)—Ce—O(3')	136.3 (2)	O(1')—Na—O(3 ^{ix})	114.00 (10)
O(3)—Ce—O(4 ^{iv})	68.6 (1)	O(2 ^{viii})—Na—O(2 ^{vii})	79.2 (2)
O(3)—Ce—O(5)	118.2 (2)	O(2 ^{viii})—Na—O(3 ^x)	76.3 (1)
O(4 ^{iv})—Ce—O(4')	76.2 (2)	O(2 ^{vii})—Na—O(3 ^x)	154.2 (2)
O(4 ^{iv})—Ce—O(5)	141.1 (3)	O(3 ^x)—Na—O(3 ^{ix})	128.9 (2)

Symmetry codes: (i) $x - y, -y, \frac{4}{3} - z$; (ii) $-y, x - y - 1, z - \frac{1}{3}$; (iii) $1 - x, 1 + y - x, \frac{2}{3} - z$; (iv) $1 - y, x - y, z - \frac{1}{3}$; (v) $1 - x, y - x, \frac{2}{3} - z$; (vi) $x, y, z - 1$; (vii) $1 + y - x, -x, z - \frac{2}{3}$; (viii) $1 + y, x, 1 - z$; (ix) $y, x - 1, 1 - z$; (x) $1 + y - x, 1 - x, z - \frac{2}{3}$.

Table 3. *Hydrogen-bonding geometry* (Å, °)

No e.s.d.'s are given for H...O(acceptor) distances (H...A) since H-atom positions are fixed. The criteria for inclusion in this tabulation are (a) the H atom involved in the hydrogen bond must be closer to the O(acceptor) atom than is the O(donor) atom (D), and (b) the O(donor) to O(acceptor) distance (D...A) is < 3.2 Å.

D—H...A	H...A	D...A	D—H...A
O(5)—H...O(1 ⁱⁱⁱ)	2.70	2.926 (10)	93.7 (4)
O(5)—H...O(4 ^{iv})	2.16	2.976 (9)	140.3 (4)
O(5)—H...O(2 ⁱⁱ)	2.50	3.114 (11)	120.9 (5)
O(5)—H ⁱ ...O(1 ⁱⁱ)	2.70	2.815 (11)	86.8 (5)
O(5)—H ⁱ ...O(4 ⁱⁱⁱ)	2.16	2.992 (8)	142.8 (5)

Symmetry codes: (i) $x - y, -y, \frac{4}{3} - z$; (ii) $-y, x - y - 1, z - \frac{1}{3}$; (iii) $1 - x, 1 + y - x, \frac{2}{3} - z$; (xi) $-y, x - y, z - \frac{1}{3}$; (xii) $-x, y - x, \frac{2}{3} - z$.

Unit-cell parameters were obtained from a symmetry-constrained least-squares fit. Scan widths were $(1.50 + 0.35 \tan \theta)^\circ$ in ω and the ratio of background to scan time was 0.5. No decay correction was made but the data were corrected for Lorentz and polarization effects. An analytical absorption correction was applied using *AGNOST* (de Meulenaer & Tompa, 1965). The intensity data were consistent with Laue group $\bar{3}m1$; the observed reflection-limiting condition $(00l, l = 3n)$ was then consistent with only the enantiomeric space groups $P3_121$ and $P3_221$. It was shown for the La analog (Blackburn & Gerkin, 1994) that capping faces $\{101\}$ are associated with $P3_121$ samples while capping faces $\{\bar{1}01\}$ are associated with $P3_221$ samples. For the experimental sample the space group $P3_221$ was given initial preference on this basis. The atomic coordinates of Na, S and La from the La analog were used as starting positions for Na, S and Ce in the present study except that each z coordinate was replaced by $-z$. Fourier difference methods were used to assign the remaining atoms. When the structure was analyzed in the potential alternative space group $P3_121$, the final R factor increased by 0.97%. Moreover, in a separate calculation, *SHELXL93* (Sheldrick, 1993) was used to refine the Flack x

parameter (Flack, 1983; Bernardinelli & Flack, 1985) simultaneously with the coordinates in $P3_221$ and gave the value $x = 0.061$ (29). These results confirmed the space group and coordinate assignments given in this report, and thus determined the absolute structure.

The refinement of the water O atom, O(5), proceeded just as described for its O(5) analog in the La salt (Blackburn & Gerkin, 1994) and was therefore modelled as disordered equally over two sites 0.57 Å apart. The H atom was modelled as ordered, also as previously: since the H-atom position was not determinable from difference maps it was fixed consistent with average water molecule geometry for the same hydrate class (N) as was observed for the La analog (Blackburn & Gerkin, 1994). Accordingly, the water O—H distance is 0.97 Å, the H—O—H angle is 104.7° and the Ce—O(5)—H angles are 120.1 (1)°. The H-atom isotropic displacement parameter was taken from the neutron-diffraction powder data for the La analog (W. J. Reppart, private communication).

The maximum positive residual electron density, $1.63 \text{ e } \text{Å}^{-3}$, occurred 0.07 Å from the Na atom while the next highest, $0.56 \text{ e } \text{Å}^{-3}$, occurred 2.52 and 2.61 Å from Na and Ce, respectively. The maximum negative density, $-0.56 \text{ e } \text{Å}^{-3}$, occurred 0.44 Å from O(5).

Rigid-body analysis of the NaO polyhedron, the CeO polyhedron and the sulfate ion was performed using the program *THMA11* (Trueblood, 1986), which is based on the work of Schomaker & Trueblood (1968). Neither the NaO nor the CeO polyhedron conformed to rigid-body behavior, but the sulfate ion did. The average magnitude of the differences in the mean-square displacement amplitudes along the interatomic vectors for the ten unique atom pairs of the sulfate ion was $28(22) \times 10^{-4} \text{ Å}^2$. The Hirshfeld (1976) rigid-bond test, extended as described by Rosenfield, Trueblood & Dunitz (1978), is satisfied and the corrected S—O bond lengths are given in this report. The rigid-body-corrected S—O bond distances are as follows: S—O(1) 1.475, S—O(2) 1.488, S—O(3) 1.485, S—O(4) 1.479 Å. These values may be compared with the corresponding rigid-body values found for the La analog: 1.481, 1.484, 1.486 and 1.475 Å, respectively.

Data collection: *MSCIAFC Diffractometer Control Software* (Molecular Structure Corporation, 1988). Cell refinement: *MSCIAFC Diffractometer Control Software*. Data reduction: *TEXSAN* (Molecular Structure Corporation, 1989). Program(s) used to refine structure: *TEXSAN*. Molecular graphics: *ORTEPII* (Johnson, 1976). Software used to prepare material for publication: *TEXSAN*.

We thank Dr W. J. Reppart for stimulating our initial interest in these double sulfates and for making unpublished results available. We thank Dr J. C. Gallucci for help of various sorts. We thank Dr Trueblood for the *THMA11* program. We gratefully acknowledge the partial support of this research through the purchase of the diffractometer system with an NIH grant.

Lists of structure factors, anisotropic displacement parameters, H-atom coordinates, complete geometry and material related to the rigid-body analysis have been deposited with the IUCr (Reference: BR1102). Copies may be obtained through The Managing Editor, International Union of Crystallography, 5 Abbey Square, Chester CH1 2HU, England.

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Ta₃S_{2-x}Se_x, Partial Substitution of Sulfur for Sulfur into a Tantalum-Rich Sulfide

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Abstract

In the pentagonal antiprismatic columnar structure of Ta₃S₂, sulfur can be partially replaced by selenium. Selenium accumulates preferentially at one of two independent sites, which leads to unusually short Ch—Ch [2.973 (3) Å] contacts, where Ch denotes a chalcogen site.

Comment

The structural chemistries of tantalum-rich sulfides and selenides differ drastically. Whereas the structure of Ta₂S (Franzen & Smeggil, 1969) results from the interpenetration of centred Ta₁₃ icosahedra along a

pseudo fivefold axis of symmetry, Ta₂Se forms a unique layered-type structure which is symmetry related to b.c.c. Ta (Harbrecht, 1989). In order to gain information about the relative stability of the distinct structural arrangements, we systematically investigated the Ta-rich region of the ternary Ta—S—Se system.

Entropically stabilized high-temperature phases of Ta₅(S,Se)₂, the (Nb,Ta)₅S₂ structure type (Yao & Franzen, 1991) and Ta₉(S,Se)₄ are accessible above temperatures of approximately 1700 K. These b.c.c.-related phases have no pendants in the binary systems (Harbrecht & Degen, 1993). In contrast to earlier reports (Nanjundaswamy & Hughbanks, 1992), we found that a significant amount of selenium can be substituted for sulfur in the structures of the sulfides. Such a mixed chalcogenide of the Ta₃S₂ structure type (Wada & Onoda, 1989; Kim, Nanjundaswamy & Hughbanks, 1991) is the subject of this report.

The structure of Ta₃(S,Se)₂ consists of distorted Ta₁₃ icosahedra which interpenetrate along the *c* axis forming columns of face-shared pentagonal antiprismatic clusters, ∞¹[TaTa₁₀/2]. Very short Ta—Ta contacts exist [Ta4—Ta4 2.811 (2) Å, compared to 2.860 Å in b.c.c.—Ta] in the centre of the pentagonal antiprismatic TaTa₁₀ clusters. The columns of condensed clusters are linked to one another, parallel to *a*, by weaker Ta1—Ta3 [3.158 (1) Å] and Ta3—Ta3 [3.650 (1) Å] contacts and strong heteronuclear bonds Ch2—Ta [2.455 (2)—2.492 (2) Å], where Ch is S or Se.

As can be seen from Fig. 2, these complex layers are also covered by Ch1 over triangular faces [Ch1—Ta1 and 2 × Ch1—Ta2 2.472 (2)—2.529 (2) Å]. Heteronuclear interlayer contacts seem to be weaker [Ch1—Ta2 2.709 (2), Ch2—Ta1 2.547 (2) Å] than interactions within the two-dimensional arrays. The specific stacking of the layers along *b* results in empty channels running

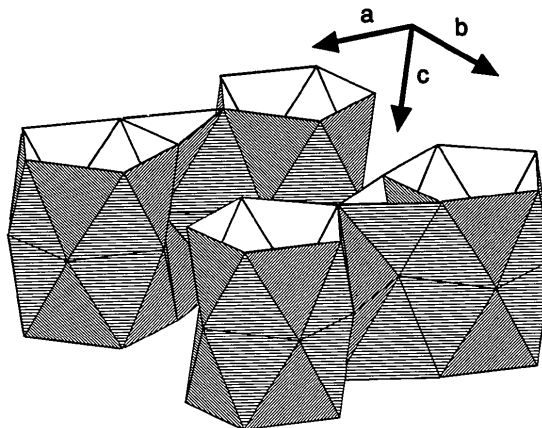


Fig. 1. Polyhedral representation of the tantalum partial structure of Ta₃(S,Se)₂ depicted as pentagonal antiprismatic columns connected via Ta₄ tetrahedra. Since the columns are centred by Ta₄ (which are omitted for clarity) the layers are tetrahedrally close packed.